

Exponential decay of mutual information for Gibbs states of local Hamiltonians

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Joint work with: **Andreas Bluhm** (U. Copenhagen)

Antonio Pérez-Hernández (UNED, Spain) .

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Granada, 24th May 2022

Andreas Bluhm

University of Copenhagen



Antonio Pérez-Hernández

Universidad Nacional de Educación a
Distancia



MOTIVATION

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Describe the **correlation properties** of **Gibbs states** of local Hamiltonians.

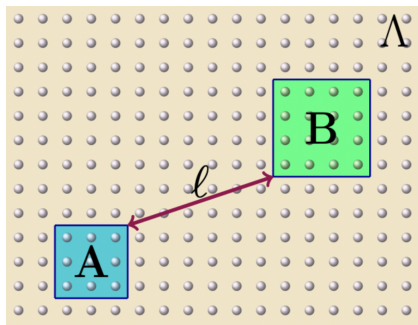
- **Hamiltonian:** $H_\Lambda = H_A + H_B + H_{(A \cup B)^c} + H_{\partial A} + H_{\partial C}$,
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Questions:

For non-commuting Hamiltonians:

$$e^{-\beta(H_A + H_B)} \approx e^{-\beta H_A} e^{-\beta H_B} ?$$

$$\text{tr}_{A^c}[\rho^\Lambda] \otimes \text{tr}_{B^c}[\rho^\Lambda] := (\rho^\Lambda)_A \otimes (\rho^\Lambda)_B \approx$$

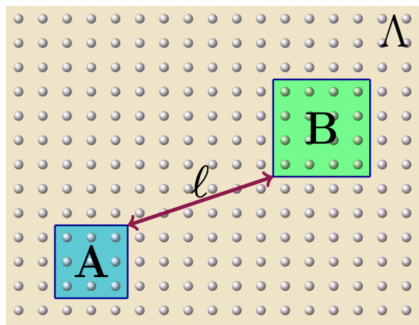
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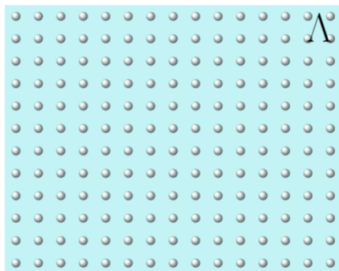
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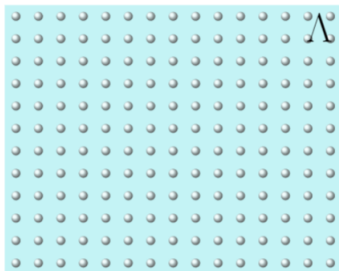
- A finite lattice $\Lambda \subset \mathbb{Z}^D$.
- For each site $x \in \Lambda$, $\mathcal{H}_x \equiv \mathbb{C}^d$.
- $\mathcal{H}_\Lambda = \bigotimes_{x \in \Lambda} \mathcal{H}_x \equiv (\mathbb{C}^d)^{\otimes |\Lambda|}$.
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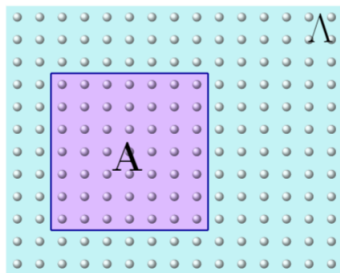
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Hamiltonian

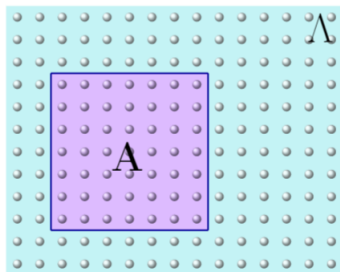
- $H_\Lambda = \sum_{X \subset \Lambda} H_X$, with $\|H_X\| \leq J$ and $\|H_X\| = 0$ if $\text{diam}(X) > r$.
- $\rho^\Lambda = \rho^\Lambda(\beta) = e^{-\beta H_\Lambda} / \text{Tr}[e^{-\beta H_\Lambda}]$.

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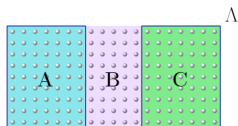


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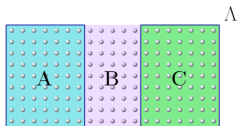
$$\Psi_X(Q) := \text{Tr}[\rho^X Q], \quad Q \in \mathfrak{A}_X$$

$$\text{Corr}_{\rho^\Lambda}(A : C) := \sup_{\|O_A\|, \|O_C\| \leq 1} |\Psi_\Lambda(O_A O_C) - \Psi_\Lambda(O_A) \Psi_\Lambda(O_C)|$$

Decay:

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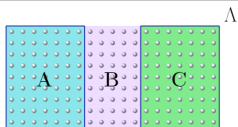
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Operator Correlation		Temperature	
		Low T	High T
Dimension	1 D	Exp. (~ Araki, '69)	
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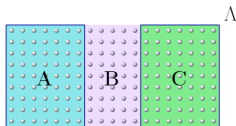
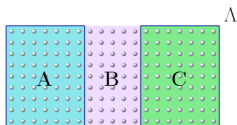
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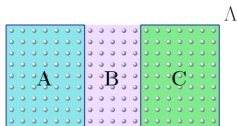
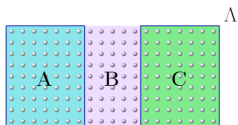
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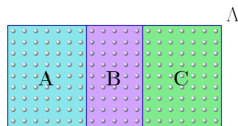
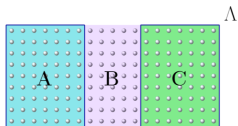
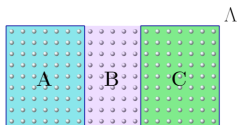
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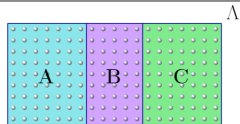
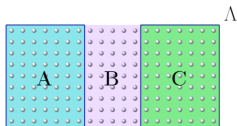
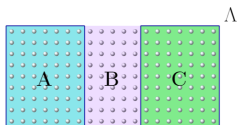
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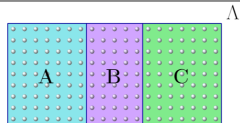
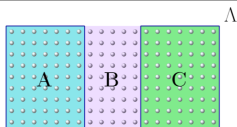
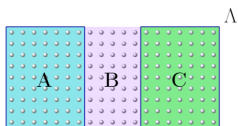
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Conditional Mutual Information		Temperature	
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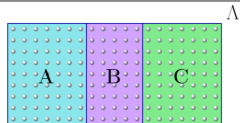
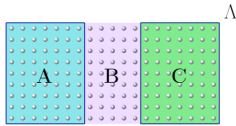
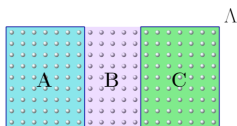
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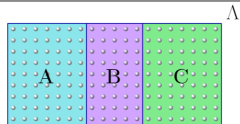
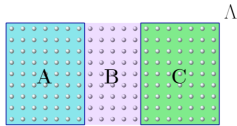
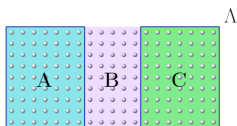
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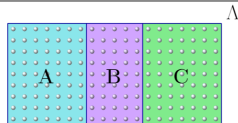
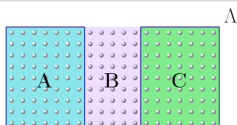
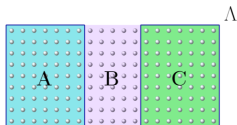
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$$I_\rho(A : C | B) := S(\rho_{AB}) + S(\rho_{BC}) - S(\rho_B) - S(\rho_{ABC})$$

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Decay:

$$\text{Corr}_{\rho^\Lambda}(A : C) \leq f(d(A : C))$$

Operator Correlation		Temperature	
		Low T	High T
Dimension	1 D	Exponential	
	Large D		

Decay:

$$I_{\rho^\Lambda}(A : C) \leq f'(d(A : C))$$

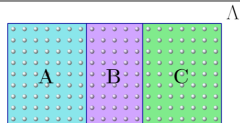
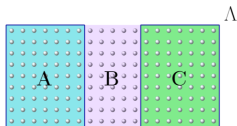
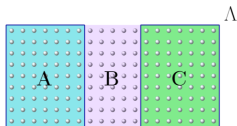
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OVERVIEW



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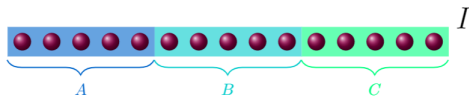
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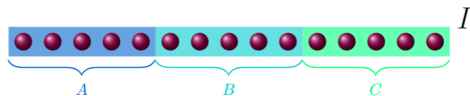
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EXPONENTIAL UNIFORM CLUSTERING

There exist $\mathcal{K}, \gamma > 0$ such that for every finite lattice $\Lambda = ABC$,

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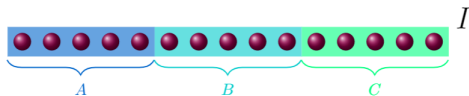
Conjecture:

In 1D, there are no thermal phase transitions.



Exponential uniform clustering holds in 1D at every $\beta > 0$.

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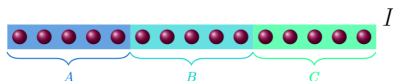
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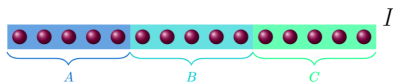
THEOREM (ARAKI, '69)

- ▶ $\mathfrak{A}_{\mathbb{Z}}$ algebra of quasi-local observables.
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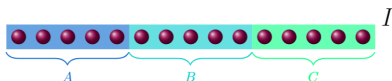
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↓ (easy) ↑ **Thm.** (Bluhm-C.-Perez Hernandez '21)

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$$I_\rho(A : B) := D(\rho_{AB} \| \rho_A \otimes \rho_B).$$

By Hölder and Pinsker's inequalities:

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Decay of mutual information \Rightarrow Decay of correlations
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There are states with small operator correlation and large mutual information in quantum data hiding (Hayden et al. '04).

AREA LAWS

$$I_\rho(A : A^c) \leq O(|\partial A|).$$

Wolf et al., '08: Area laws for the mutual information (improved: Kuwahara et al. '20).

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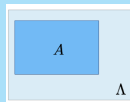
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$$\hat{D}(\rho_{AC} \parallel \sigma_{AC}) := \text{Tr} \left[\rho_{AC} \log \left(\rho_{AC}^{1/2} \sigma_{AC}^{-1} \rho_{AC}^{1/2} \right) \right].$$

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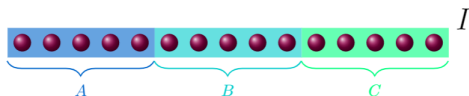
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Assuming exponential uniform clustering, there exist $\mathcal{K}, \gamma > 0$ such that for every finite interval $I = ABC$ and $\rho = \rho^I = e^{-\beta H_I} / \text{Tr}[e^{-\beta H_I}]$,

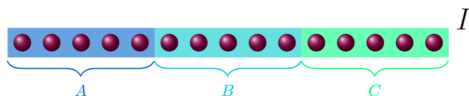
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The following are equivalent:

- 1 Exponential decay of correlations for the infinite-chain thermal states.
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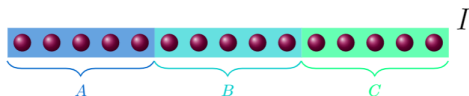
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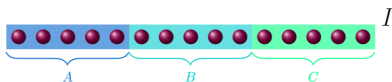
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SKETCH OF THE PROOF



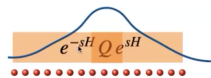
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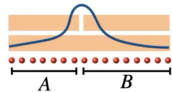
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Ingredients:

- Imaginary-time **Lieb-Robinson bounds** (Araki, '69)



- Araki's **expansionals** (Araki, '69): $E = e^{H_A + H_B} e^{-H_{AB}}$



- Exponential uniform clustering**

- Local indistinguishability** (Brandao-Kast., '19): $(\rho^{ABC})_A \approx (\rho^{AB})_A$



CONDITIONAL MUTUAL INFORMATION

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STRONG SUBADDITIVITY AND RECOVERABILITY (Lieb-Ruskai, '73, Petz, '86)

$$I_\rho(A : C|B) \geq 0$$

and $I_\rho(A : C|B) = 0$ if, and only if, ρ_{ABC} is a quantum Markov chain ($A \leftrightarrow B \leftrightarrow C$):

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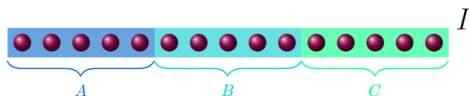
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APPROXIMATE FACTORIZATION OF GIBBS STATES



BS - CONDITIONAL MUTUAL INFORMATION

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For two positive states $\rho_{ABC}, \sigma_{ABC} \in \mathcal{H}_{ABC}$ such that $\sigma_{ABC} = \rho_{AB} \otimes \mathbb{1}_C / d_C$ and a $\mathcal{T} := \mathbb{1}_A / d_A \otimes \text{Tr}_A$,

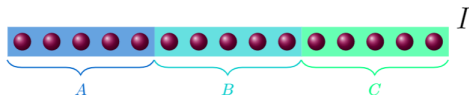
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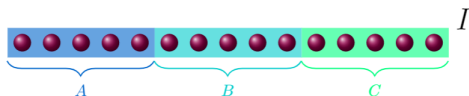
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APPLICATIONS TO MLSI



EXPONENTIAL DECAY MIXING CONDITION

Assuming exponential uniform clustering, there exist $\mathcal{K}, \gamma > 0$ such that for every finite interval $I = ABC$ and $\rho = \rho^I = e^{-\beta H_I} / \text{Tr}[e^{-\beta H_I}]$,

$$\|\rho_A^{-1} \otimes \rho_C^{-1} \rho_{AC} - \mathbb{1}_{AC}\| \leq \mathcal{K} e^{-\gamma|B|}$$

MLSI FOR HEAT-BATH DYNAMICS (Bardet-C.-Lucia-Perez Garcia-Rouze, '21)

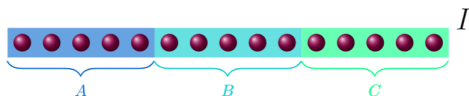
Theorem 7. Let $\Lambda \subset \mathbb{Z}$ be a finite chain. Let $\Phi : \Lambda \rightarrow \mathcal{A}_\Lambda$ be a k -local commuting potential and $H_\Lambda = \sum_{x \in \Lambda} \Phi(x)$ be its corresponding Hamiltonian, and denote by σ_Λ its Gibbs state. Let \mathcal{L}_Λ^* be the generator of the heat-bath dynamics. Then, if Assumptions 1 and 2 hold, the MLSI constant of \mathcal{L}_Λ^* is strictly positive and independent of $|\Lambda|$.

Assumption 1 (mixing condition). Let $\Lambda \subset \mathbb{Z}$ be a finite chain, and let $C, D \subset \Lambda$ be the union of non-overlapping finite-sized segments of Λ . Let σ_Λ be the Gibbs state of a commuting Hamiltonian. The following inequality holds for certain positive constants K_1, K_2 independent of Λ, C, D :

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where $d(C, D)$ is the distance between C and D , i.e., the minimum distance between two segments of C and D .

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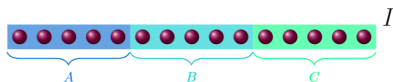
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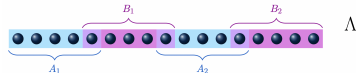
APPLICATIONS TO MLSI



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For translation-invariant interactions, there exist $\mathcal{K}, \gamma > 0$ such that for every finite interval $I = ABC$ and $\rho = \rho^I = e^{-\beta H_I} / \text{Tr}[e^{-\beta H_I}]$,

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MLSI FOR HEAT-BATH DYNAMICS (Bardet-C.-Gao-Lucia-Perez Garcia-Rouze, '21)

Positive MLSI ($O(\log(|I|))$) for Davies generators in 1D at any $\beta > 0$.

Theorem 3.1. Let $\Lambda = [1, n]$. For any $\beta > 0$, we denote by $\sigma \equiv \sigma^\beta$ the Gibbs state of a finite-range, translation-invariant, commuting Hamiltonian at inverse temperature $\beta > 0$. Consider \mathcal{L}_Λ^D the Davies generator of a quantum Markov semigroup $\{e^{t\mathcal{L}_\Lambda^D}\}_{t \geq 0}$ with unique fixed point σ . Then, there exists $\alpha_n = \Omega(\ln(n)^{-1})$ such that, for all $\rho \in \mathcal{D}(\mathcal{H}_\Lambda)$ and all $t \geq 0$,

$$D(\rho_t \| \sigma) \leq e^{-\alpha_n t} D(\rho \| \sigma), \quad (25)$$

where $\rho_t := e^{t\mathcal{L}_\Lambda^D}(\rho)$. Moreover, $\alpha_n = e^{-O(\beta)}$ as a function of β .

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There exists $\mathcal{G}, \mathcal{K} > 1$ such that for every finite interval $I \subset \mathbb{Z}$, split into $I = A_1 A_2 \dots A_n$ with $|A_j| = m$ for every $j = 1, \dots, n$, and for $\rho = \rho^I$ the Gibbs state on I ,

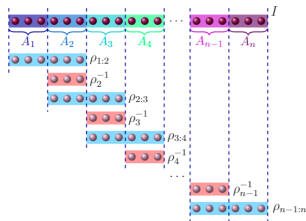
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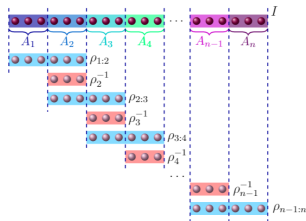
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