# Entropy decay for Davies semigroups of a one dimensional quantum lattice

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#### PROBLEM

Velocity of convergence of certain quantum dissipative evolutions to their thermal equilibriums.

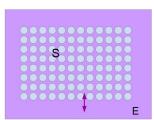
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 $\Rightarrow$  Open quantum many-body systems.

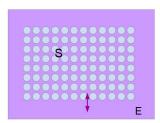


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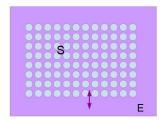
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A quantum Markov semigroup is a 1-parameter continuous semigroup  $\{\mathcal{T}_t^*\}_{t\geq 0}$  of completely positive, trace preserving (CPTP) maps (a.k.a. quantum channels) in  $\mathcal{S}_{\Lambda}$ .

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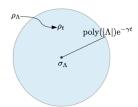
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We say that  $\mathcal{L}^*_{\Lambda}$  satisfies **rapid mixing** if

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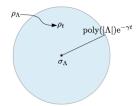
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The MLSI constant of  $\mathcal{L}^*_{\Lambda}$  is defined as:

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Using the spectral gap (Kastoryano-Temme '13):

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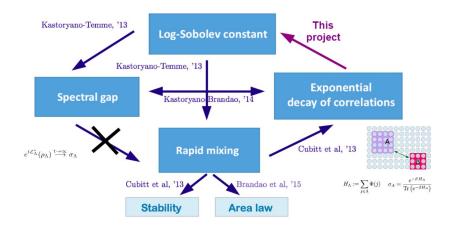
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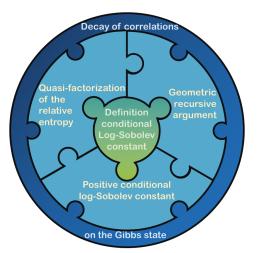
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#### STRATEGY

Used in (C.-Lucia-Pérez García '18) and (Bardet-C.-Lucia-Pérez García-Rouzé, '19).



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The **conditional MLSI constant** of  $\mathcal{L}^*_{\Lambda}$  on  $A \subset \Lambda$  is defined by

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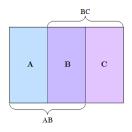
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## Quasi-factorization of the relative entropy



## QUASI-FACTORIZATION OF THE RELATIVE ENTROPY

Given  $\Lambda = ABC$ , it is an inequality of the form:

$$D(\rho_{\Lambda} \| \sigma_{\Lambda}) \leq \xi(\sigma_{ABC}) \left[ D_{AB}(\rho_{\Lambda} \| \sigma_{\Lambda}) + D_{BC}(\rho_{\Lambda} \| \sigma_{\Lambda}) \right] ,$$

for  $\rho_{\Lambda}, \sigma_{\Lambda} \in \mathcal{D}(\mathcal{H}_{ABC})$ , where  $\xi(\sigma_{ABC})$  depends only on  $\sigma_{ABC}$  and measures how far  $\sigma_{AC}$  is from  $\sigma_{A} \otimes \sigma_{C}$ .

## Example: Tensor product fixed point

(C.-Lucia-Pérez García '18) 
$$\mathcal{L}_{\Lambda}^*(\rho_{\Lambda}) = \sum_{x \in \Lambda} (\sigma_x \otimes \rho_{x^c} - \rho_{\Lambda})$$
 
$$D_x(\rho_{\Lambda} \| \sigma_{\Lambda}) := D(\rho_{\Lambda} \| \sigma_{\Lambda}) - D(\rho_{x^c} \| \sigma_{x^c})$$



$$\sigma_{\Lambda} = \bigotimes_{x \in \Lambda} \sigma_x,$$



$$D(\rho_{\Lambda}||\sigma_{\Lambda}) \leq$$



$$\leq \sum_{x \in \Lambda} D_x(\rho_{\Lambda}||\sigma_{\Lambda})$$

$$\alpha_{\Lambda}(\mathcal{L}_x^*) := \inf_{\rho_{\Lambda} \in S_{\Lambda}} \frac{-\operatorname{tr}[\mathcal{L}_x^*(\rho_{\Lambda})(\log \rho_{\Lambda} - \log \sigma_{\Lambda})]}{2D_x(\rho_{\Lambda}||\sigma_{\Lambda})}$$

$$\frac{\sum_{\mathbf{x} \in \Lambda} \frac{\operatorname{tr}[\mathcal{L}(\mathbf{x}) | \operatorname{the}_{\mathbf{x}} - \operatorname{log}_{\mathbf{x}}]}{2 \mathcal{H}(\mathbf{x}) | \mathbf{x}|}}{2 \mathcal{H}(\mathbf{x}) | \mathbf{x}|} \leq \sum_{x \in \Lambda} \frac{-\operatorname{tr}[\mathcal{L}_{x}^{*}(\rho_{\Lambda}) (\log \rho_{\Lambda} - \log \sigma_{\Lambda})]}{2 \alpha_{\Lambda}(\mathcal{L}_{x}^{*})}$$

$$\leq \frac{1}{2\inf_{x\in\Lambda}\alpha_{\Lambda}(\mathcal{L}_{x}^{*})}\sum_{x\in\Lambda} -\operatorname{tr}[\mathcal{L}_{x}^{*}(\rho_{\Lambda})(\log\rho_{\Lambda} - \log\sigma_{\Lambda})]$$

$$= \frac{1}{2\inf \alpha_{\Lambda}(\mathcal{L}_{x}^{*})} \left( -\operatorname{tr}[\mathcal{L}_{\Lambda}^{*}(\rho_{\Lambda})(\log \rho_{\Lambda} - \log \sigma_{\Lambda})] \right)$$

$$\leq \left(-\operatorname{tr}[\mathcal{L}_{\Lambda}^*(\rho_{\Lambda})(\log \rho_{\Lambda} - \log \sigma_{\Lambda})]\right).$$

Let  $\sigma_{\Lambda} = \frac{e^{-\beta H_{\Lambda}}}{\text{tr}[e^{-\beta H_{\Lambda}}]}$  be the Gibbs state of finite-range, commuting Hamiltonian.

#### HEAT-BATH GENERATOR

The heat-bath generator is defined as:

$$\mathcal{L}_{\Lambda}^{H;*}(\rho_{\Lambda}) := \sum_{x \in \Lambda} \left( \sigma_{\Lambda}^{1/2} \sigma_{x^c}^{-1/2} \rho_{x^c} \sigma_{x^c}^{-1/2} \sigma_{\Lambda}^{1/2} - \rho_{\Lambda} \right)$$

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The Davies generator is given by:

$$\mathcal{L}_{\Lambda}^{D}(X) := i[H_{\Lambda}, X] + \sum_{x \in \Lambda} \mathcal{L}_{x}^{D}(X)$$

where the  $\mathcal{L}_x^D$  are defined in terms of the Fourier coefficients of the correlation functions in the bath and the ones of the system couplings.

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The **Schmidt generator** (Bravyi-Vyalyi '05) can be written as:

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## SPECTRAL GAP FOR DAVIES AND HEAT-BATH (Kastoryano-Brandao, '16)

Let  $\mathcal{L}_{\Lambda}^{H,D;*}$  be the **heat-bath** or **Davies** generator in 1D. Then,  $\mathcal{L}_{\Lambda}^{H,D;*}$  has a positive spectral gap that is independent of the system size, for every temperature

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#### Assumption 1

In a tripartite Hilbert space  $\mathcal{H}_A \otimes \mathcal{H}_C \otimes \mathcal{H}_B$ , A and B not connected, we have

$$\left\| \sigma_A^{-1/2} \otimes \sigma_B^{-1/2} \sigma_{AB} \sigma_A^{-1/2} \otimes \sigma_B^{-1/2} - \mathbb{1}_{AB} \right\|_{\infty} \le K < \frac{1}{2}.$$

In particular, Gibbs states at high enough temperature satisfy this.

#### Assumption 2

For any  $B \subset \Lambda$ ,  $B = B_1 \cup B_2$ , it holds:

$$D_B(\rho_{\Lambda}||\sigma_{\Lambda}) \le f(\sigma_{B\partial}) \left( D_{B_1}(\rho_{\Lambda}||\sigma_{\Lambda}) + D_{B_2}(\rho_{\Lambda}||\sigma_{\Lambda}) \right)$$

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In 1D, if Assumptions 1 and 2 hold, for a k-local commuting Hamiltonian, the **heat-bath** dynamics has a positive log-Sobolev constant.



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- $\bullet$   $H_{\Lambda}$  is classical.
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- $\bullet$   $\Lambda$  is 1D.

Then, there exists a local quantum Markov semigroup with fixed point  $\sigma_{\Lambda}$ , the Gibbs state of  $H_{\Lambda}$ , such that it has a positive **MLSI constant** which is independent of the system size.

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## Main result

## MLSI FOR 1D DAVIES GENERATORS, (Bardet-C.-Gao-Lucia-Pérez García-Rouzé, '21)

Let  $\mathcal{L}_{\Lambda}^{D;*}$  be a **Davies** generator with unique fixed point  $\sigma_{\Lambda}$  given by the Gibbs state of a commuting, finite-range, translation-invariant Hamiltonian at any temperature in 1D. Then,  $\mathcal{L}_{\Lambda}^{D;*}$  satisfies a positive MLSI  $\alpha(\mathcal{L}_{\Lambda}^{D;*}) = \Omega(\ln(|\Lambda|)^{-1})$ .

## Rapid mixing:

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# Proof: Conditional relative entropies + Quasi-factorization





Conditional relative entropies:  $D_A(\rho_{\Lambda} \| \sigma_{\Lambda}) := D(\rho_{\Lambda} \| \sigma_{\Lambda}) - D(\rho_{A^c} \| \sigma_{A^c})$ ,  $D_A^E(\rho_{\Lambda} \| \sigma_{\Lambda}) := D(\rho_{\Lambda} \| E_A^*(\rho_{\Lambda}))$ .

## Quasi-factorization (C.-Lucia-Pérez García '18

Let  $\mathcal{H}_{ABC}$  and  $\rho_{ABC}, \sigma_{ABC} \in \mathcal{S}_{ABC}$ . The following holds

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$$\xi(\sigma_{AC}) = \frac{1}{1 - 2\|\sigma_A^{-1/2} \otimes \sigma_C^{-1/2} \sigma_{AC} \sigma_A^{-1/2} \otimes \sigma_C^{-1/2} - \mathbb{1}_{AC}\|}.$$

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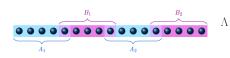
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# PROOF: QUASI-FACTORIZATION





 $\sigma_{\Lambda} = \frac{e^{-\beta H_{\Lambda}}}{\operatorname{tr}(e^{-\beta H_{\Lambda}})}$  is the Gibbs state of a k-local, commuting Hamiltonian  $H_{\Lambda}$ .

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Let  $A \cup B = \Lambda \subset \mathbb{Z}$  and  $\rho_{\Lambda}, \sigma_{\Lambda} \in \mathcal{S}_{\Lambda}$ . The following holds

$$D(\rho_{\Lambda}||\sigma_{\Lambda}) \leq \xi(\sigma_{A^c B^c}) \left[ D_A(\rho_{\Lambda}||\sigma_{\Lambda}) + D_B(\rho_{\Lambda}||\sigma_{\Lambda}) \right],$$

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# QUASI-FACTORIZATION FOR QUANTUM MARKOV CHAINS (Bardet-C.-Lucia-Pérez García-Rouzé'19

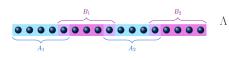
Since  $\sigma_{\Lambda}$  is a QMC between  $A_i \leftrightarrow \partial (A_i) \leftrightarrow (A_i \cup \partial A_i)^c$ , then

$$D_A(\rho_{\Lambda}||\sigma_{\Lambda}) \le \sum_i D_{A_i}(\rho_{\Lambda}||\sigma_{\Lambda})$$

$$\sigma_{\Lambda} = \bigoplus_{j \in J} \sigma_{A_i(\partial a_i)_j^L} \otimes \sigma_{(\partial a_i)_j^R(A_i \cup \partial A_i)^C}$$



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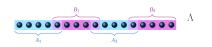
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## Proof: Decay of Correlations





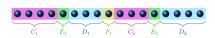
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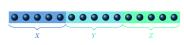
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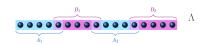


DECAY OF CORRELATIONS, (Bluhm-C.-Pérez Hernández, '21)

Let  $\sigma_{XYZ}$  be the Gibbs state of a finite-range, translation-invariant Hamiltonian. There is  $\ell \mapsto \delta(\ell)$  with exponential decay such that:

$$\left\|\sigma_X^{-1} \otimes \sigma_Z^{-1} \sigma_{XZ} - \mathbb{1}_{XZ}\right\|_{\infty} \le \delta(|Y|).$$

# PROOF: DECAY OF CORRELATIONS





### Quasi-factorization

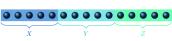
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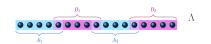
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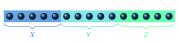
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# Proof: Geometric recursive argument





Let us recall:  $D_A(\rho_{\Lambda} || \sigma_{\Lambda}) := D(\rho_{\Lambda} || \sigma_{\Lambda}) - D(\rho_{A^c} || \sigma_{A^c})$ ,  $D_{\Lambda}^{E}(\rho_{\Lambda} \| \sigma_{\Lambda}) := D(\rho_{\Lambda} \| E_{\Lambda}^{*}(\rho_{\Lambda}))$ .

# Comparison between conditional relative entropies (Bardet-C.-Rouzé, '20)

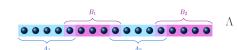
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## Proof: Positive CMLSI



# REDUCTION OF CONDITIONAL RELATIVE ENTROPIES (Gao-Rouzé, '21)

$$D(\rho_{\Lambda} || E_{A_i}^*(\rho_{\Lambda})) \le 4k_{A_i} \sum_{j \in A_i} D(\rho_{\Lambda} || E_j^*(\rho_{\Lambda}))$$

### REDUCTION FROM CMLSI TO GAP

$$k_{A_i} \propto \frac{1}{\ln \lambda}$$

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The CMLSI of the local generators is positive:

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The conditional expectations associated to Davies and heat-bath dynamics coincide.

### Conclusion

For  $\mathcal{L}_{\Lambda}^{D;*}$ , there is a positive MLSI constant  $\alpha(\mathcal{L}_{\Lambda}^{D;*}) = \Omega(\ln |\Lambda|^{-1})$ . Therefore,  $\mathcal{L}_{\Lambda}^{D;*}$  has rapid mixing.

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