

Entropy decay for 1D quantum Gibbs samplers

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Joint work with Ivan Bardet (INRIA, Paris), Nilanjana Datta (U. Cambridge), Angelo Lucia (Caltech), Cambyse Rouzé (T. U. München) and David Pérez-García (U. Complutense de Madrid).

Caltech, IQI Seminar, 16th April 2019

BASED ON:

- ① A. Capel, A. Lucia and D. Pérez-García, **Superadditivity of Quantum Relative Entropy for General States**, *IEEE Trans. on Inf. Theory*, 64 (7) (2018), 4758–4765.

- ② A. Capel, A. Lucia and D. Pérez-García, **Quantum Conditional Relative Entropy and Quasi-Factorization of the Relative Entropy**, *J. Phys. A: Math. Theor.*, 51 (2018), 484001.

- ③ I. Bardet, A. Capel, N. Datta, A. Lucia, D. Pérez-García and C. Rouzé, **Entropy decay for 1D quantum Gibbs samplers**, in preparation.

Q. information theory \longleftrightarrow **Q. many-body physics**

Communication channels \longleftrightarrow Physical interactions

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FIELD OF STUDY

Dissipative evolutions of quantum many-body systems

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Velocity of convergence of certain quantum dissipative evolutions to their thermal equilibriums.

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Provide sufficient static conditions on a Gibbs state which imply the existence of a positive log-Sobolev constant.

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- 1 QUANTUM DISSIPATIVE SYSTEMS
- 2 QUASI-FACTORIZATION OF THE RELATIVE ENTROPY
 - CONDITIONAL RELATIVE ENTROPY
 - QUASI-FACTORIZATION OF THE RELATIVE ENTROPY
- 3 LOG-SOBOLEV CONSTANT

1. QUANTUM DISSIPATIVE SYSTEMS

OPEN QUANTUM SYSTEMS

No experiment can be executed at zero temperature or be completely shielded from noise.

⇒ Open quantum many-body systems.

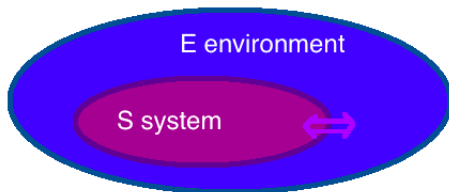


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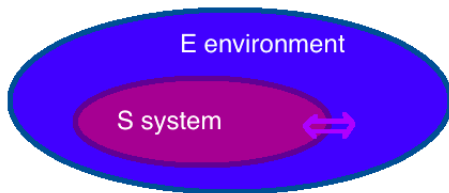


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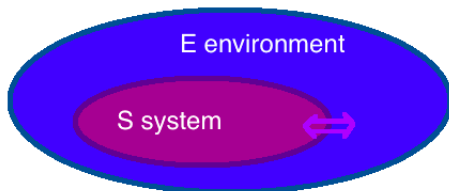


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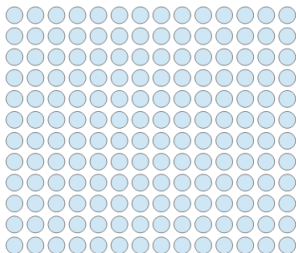


Figure: A quantum spin lattice system.

- Finite lattice $\Lambda \subset \mathbb{Z}^d$.
- To every site $x \in \Lambda$ we associate $\mathcal{H}_x (= \mathbb{C}^D)$.
- The global Hilbert space associated to Λ is $\mathcal{H}_\Lambda = \bigotimes_{x \in \Lambda} \mathcal{H}_x$.
- The set of bounded linear endomorphisms on \mathcal{H}_Λ is denoted by $\mathcal{B}_\Lambda := \mathcal{B}(\mathcal{H}_\Lambda)$.
- The set of density matrices is denoted by $\mathcal{S}_\Lambda := \mathcal{S}(\mathcal{H}_\Lambda) = \{\rho_\Lambda \in \mathcal{B}_\Lambda : \rho_\Lambda \geq 0 \text{ and } \text{tr}[\rho_\Lambda] = 1\}$.

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Physical evolution: $\rho \mapsto U \rho U^* \rightsquigarrow$ Reversible

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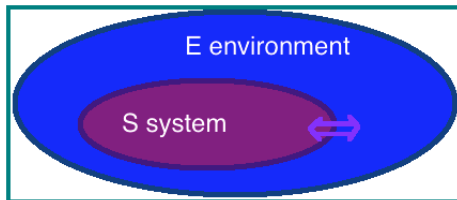


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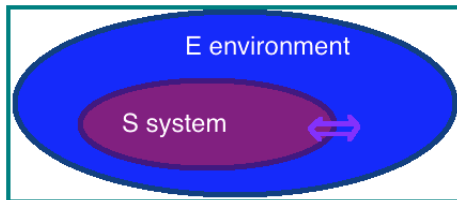


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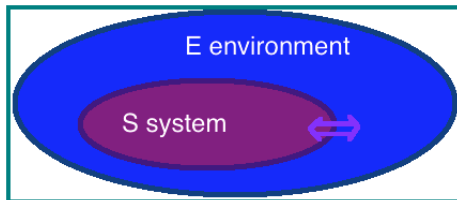


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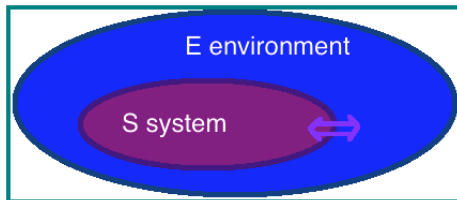


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Semigroup:

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The infinitesimal generator \mathcal{L}_Λ^* of the previous semigroup of quantum channels is usually called **Liouvillian**, or **Lindbladian**.

$$\mathcal{T}_t^* = e^{t\mathcal{L}_\Lambda^*} \Leftrightarrow \mathcal{L}_\Lambda^* = \left. \frac{d}{dt} \mathcal{T}_t^* \right|_{t=0}.$$

Notation: $\rho_t := \mathcal{T}_t^*(\rho)$.

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Recent change of perspective \Rightarrow Resource to exploit

New area:

Quantum dissipative engineering,

to create artificial evolutions in which the dissipative process works in favor (protecting the system from noisy evolutions).

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We define the **mixing time** of $\{\mathcal{T}_t^*\}$ by

$$\tau(\varepsilon) = \min \left\{ t > 0 : \sup_{\rho_\Lambda \in \mathcal{S}_\Lambda} \|\mathcal{T}_t^*(\rho) - \mathcal{T}_\infty^*(\rho)\|_1 \leq \varepsilon \right\}.$$

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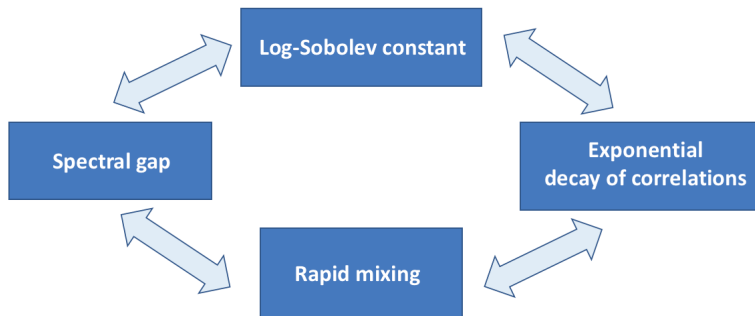
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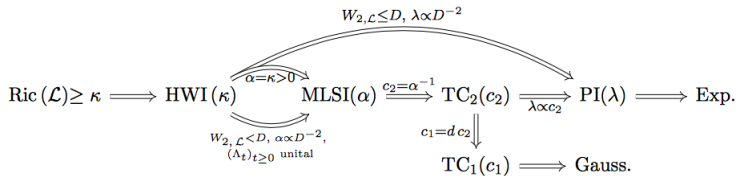
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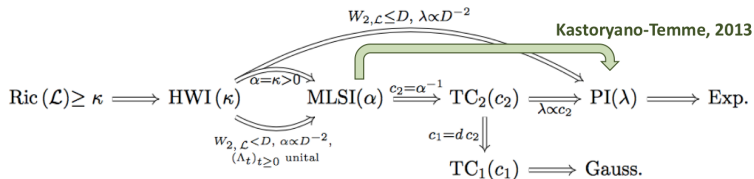
CLASSICAL SPIN SYSTEMS



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(1) Quasi-factorization of the entropy (in terms of a conditional entropy).

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Lower bound for the global log-Sobolev constant in terms of the log-Sobolev constant of a size-fixed region.

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(Cesi, Dai Pra-Paganoni-Posta, '02)

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CONDITIONAL LOG-SOBOLEV CONSTANT

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Let $\mathcal{L}_\Lambda^* : \mathcal{S}_\Lambda \rightarrow \mathcal{S}_\Lambda$ be a primitive reversible Lindbladian with stationary state σ_Λ , $A \subseteq \Lambda$. We define the **conditional log-Sobolev constant** of \mathcal{L}_Λ^* on A by

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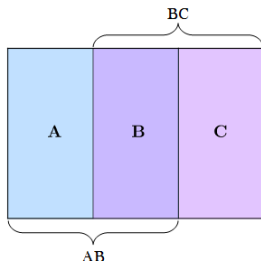
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2. QUASI-FACTORIZATION OF THE RELATIVE ENTROPY

STATEMENT OF THE PROBLEM



PROBLEM

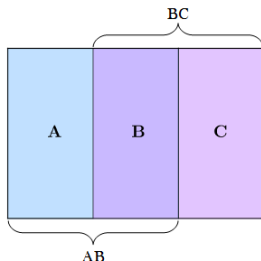
Let $\mathcal{H}_{ABC} = \mathcal{H}_A \otimes \mathcal{H}_B \otimes \mathcal{H}_C$ and $\rho_{ABC}, \sigma_{ABC} \in \mathcal{S}_{ABC}$. Can we prove something like

$$D(\rho_{ABC} || \sigma_{ABC}) \leq \xi(\sigma_{ABC}) [D_{AB}(\rho_{ABC} || \sigma_{ABC}) + D_{BC}(\rho_{ABC} || \sigma_{ABC})] ?$$

QUANTUM RELATIVE ENTROPY

$$D(\rho || \sigma) = \text{tr} [\rho (\log \rho - \log \sigma)]$$

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CLASSICAL CASE, Dai Pra et al. '02

$$\text{Ent}_\mu(f) \leq \frac{1}{1 - 4\|h - 1\|_\infty} \mu [\text{Ent}_\mu(f | \mathcal{F}_1) + \text{Ent}_\mu(f | \mathcal{F}_2)],$$

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CLASSICAL ENTROPY AND CONDITIONAL ENTROPY

Entropy:

$$\text{Ent}_\mu(f) = \mu(f \log f) - \mu(f) \log \mu(f).$$

Conditional entropy:

$$\text{Ent}_\mu(f | \mathcal{G}) = \mu(f \log f | \mathcal{G}) - \mu(f | \mathcal{G}) \log \mu(f | \mathcal{G}).$$

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RELATIVE ENTROPY

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Let $\rho_\Lambda, \sigma_\Lambda \in \mathcal{S}_\Lambda$. The **quantum relative entropy** of ρ_Λ and σ_Λ is defined by:

$$D(\rho_\Lambda || \sigma_\Lambda) = \text{tr} [\rho_\Lambda (\log \rho_\Lambda - \log \sigma_\Lambda)].$$

PROPERTIES OF THE RELATIVE ENTROPY

Let $\mathcal{H}_{AB} = \mathcal{H}_A \otimes \mathcal{H}_B$ and $\rho_{AB}, \sigma_{AB} \in \mathcal{S}_{AB}$. The following properties hold:

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CHARACTERIZATION OF THE RE, Wilming et al. '17, Matsumoto '10

If $f : \mathcal{S}_{AB} \times \mathcal{S}_{AB} \rightarrow \mathbb{R}_0^+$ satisfies 1 – 4, then f is the relative entropy.

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CONDITIONAL RELATIVE ENTROPY

CONDITIONAL RELATIVE ENTROPY

Let $\mathcal{H}_{AB} = \mathcal{H}_A \otimes \mathcal{H}_B$. We define a **conditional relative entropy** in A as a function

$$D_A(\cdot || \cdot) : \mathcal{S}_{AB} \times \mathcal{S}_{AB} \rightarrow \mathbb{R}_0^+$$

verifying the following properties for every $\rho_{AB}, \sigma_{AB} \in \mathcal{S}_{AB}$:

① **Continuity:** The map $\rho_{AB} \mapsto D_A(\rho_{AB} || \sigma_{AB})$ is continuous.

② **Non-negativity:** $D_A(\rho_{AB} || \sigma_{AB}) \geq 0$ and

$$(2.1) \quad D_A(\rho_{AB} || \sigma_{AB}) = 0 \text{ if, and only if, } \rho_{AB} = \sigma_{AB}^{1/2} \sigma_B^{-1/2} \rho_B \sigma_B^{-1/2} \sigma_{AB}^{1/2}.$$

③ **Semi-superadditivity:** $D_A(\rho_{AB} || \sigma_A \otimes \sigma_B) \geq D(\rho_A || \sigma_A)$ and

$$(3.1) \quad \text{Semi-additivity: if } \rho_{AB} = \rho_A \otimes \rho_B, \\ D_A(\rho_A \otimes \rho_B || \sigma_A \otimes \sigma_B) = D(\rho_A || \sigma_A).$$

④ **Semi-monotonicity:** For every quantum channel \mathcal{T} ,

$$D_A(\mathcal{T}(\rho_{AB}) || \mathcal{T}(\sigma_{AB})) + D_B((\text{tr}_A \circ \mathcal{T})(\rho_{AB}) || (\text{tr}_A \circ \mathcal{T})(\sigma_{AB})) \\ \leq D_A(\rho_{AB} || \sigma_{AB}) + D_B(\text{tr}_A(\rho_{AB}) || \text{tr}_A(\sigma_{AB})).$$

REMARK

Consider for every $\rho_{AB}, \sigma_{AB} \in \mathcal{S}_{AB}$

$$D_{A,B}^+(\rho_{AB}||\sigma_{AB}) = D_A(\rho_{AB}||\sigma_{AB}) + D_B(\rho_{AB}||\sigma_{AB}).$$

Then, $D_{A,B}^+$ verifies the following properties:

- ① **Continuity:** $\rho_{AB} \mapsto D_{A,B}^+(\rho_{AB}||\sigma_{AB})$ is continuous.
- ② **Additivity:** $D_{A,B}^+(\rho_A \otimes \rho_B || \sigma_A \otimes \sigma_B) = D(\rho_A || \sigma_A) + D(\rho_B || \sigma_B)$.
- ③ **Superadditivity:** $D_{A,B}^+(\rho_{AB} || \sigma_A \otimes \sigma_B) \geq D(\rho_A || \sigma_A) + D(\rho_B || \sigma_B)$.

However, it does not satisfy the property of monotonicity.

AXIOMATIC CHARACTERIZATION OF THE CRE (C-Lucia-Pérez García, '18)

The only possible conditional relative entropy is given by:

$$D_A(\rho_{AB}||\sigma_{AB}) = D(\rho_{AB}||\sigma_{AB}) - D(\rho_B||\sigma_B)$$

for every $\rho_{AB}, \sigma_{AB} \in \mathcal{S}_{AB}$.

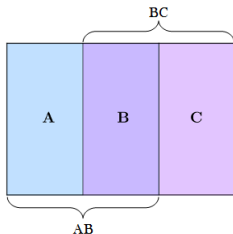


Figure: Choice of indices in $\mathcal{H}_{ABC} = \mathcal{H}_A \otimes \mathcal{H}_B \otimes \mathcal{H}_C$.

Result of **quasi-factorization** of the relative entropy, for every $\rho_{ABC}, \sigma_{ABC} \in \mathcal{S}_{ABC}$:

$$D(\rho_{ABC} || \sigma_{ABC}) \leq \xi(\sigma_{ABC}) [D_{AB}(\rho_{ABC} || \sigma_{ABC}) + D_{BC}(\rho_{ABC} || \sigma_{ABC})],$$

where $\xi(\sigma_{ABC})$ depends only on σ_{ABC} and measures how far σ_{AC} is from $\sigma_A \otimes \sigma_C$.

QUASI-FACTORIZATION FOR THE CRE (C-Lucía-Pérez García, '18)

Let $\mathcal{H}_{ABC} = \mathcal{H}_A \otimes \mathcal{H}_B \otimes \mathcal{H}_C$ and $\rho_{ABC}, \sigma_{ABC} \in \mathcal{S}_{ABC}$. Then, the following inequality holds

$$D(\rho_{ABC} || \sigma_{ABC}) \leq \frac{1}{1 - 2\|H(\sigma_{AC})\|_\infty} [D_{AB}(\rho_{ABC} || \sigma_{ABC}) + D_{BC}(\rho_{ABC} || \sigma_{ABC})],$$

where

$$H(\sigma_{AC}) = \sigma_A^{-1/2} \otimes \sigma_C^{-1/2} \sigma_{AC} \sigma_A^{-1/2} \otimes \sigma_C^{-1/2} - \mathbb{1}_{AC}.$$

Note that $H(\sigma_{AC}) = 0$ if σ_{AC} is a tensor product between A and C .

$$\begin{aligned}
 (1 - 2\|H(\sigma_{AC})\|_\infty)D(\rho_{ABC}\|\sigma_{ABC}) &\leq \\
 D_{AB}(\rho_{ABC}\|\sigma_{ABC}) + D_{BC}(\rho_{ABC}\|\sigma_{ABC}) &= \\
 = 2D(\rho_{ABC}\|\sigma_{ABC}) - D(\rho_C\|\sigma_C) - D(\rho_A\|\sigma_A). &
 \end{aligned}$$

\Leftrightarrow

$$(1 + 2\|H(\sigma_{AC})\|_\infty)D(\rho_{ABC}\|\sigma_{ABC}) \geq D(\rho_A\|\sigma_A) + D(\rho_C\|\sigma_C).$$

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This result is equivalent to:

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Recall:

- **Superadditivity.** $D(\rho_{AB}||\sigma_A \otimes \sigma_B) \geq D(\rho_A||\sigma_A) + D(\rho_B||\sigma_B)$.

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we have

$$2D(\rho_{AB}||\sigma_{AB}) \geq D(\rho_A||\sigma_A) + D(\rho_B||\sigma_B).$$

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RELATION WITH THE CLASSICAL CASE

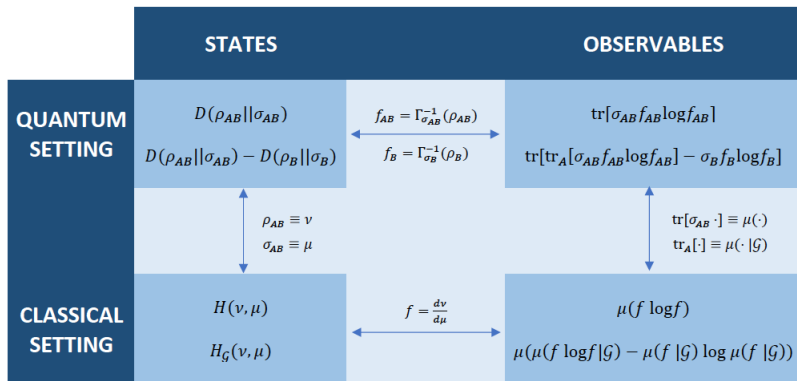


Figure: Identification between classical and quantum quantities when the states considered are classical.

3. LOG-SOBOLEV CONSTANT

QUANTUM SPIN LATTICES

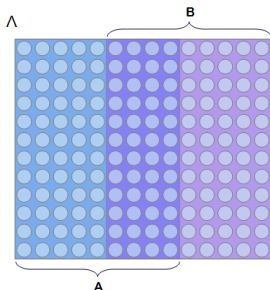


Figure: A quantum spin lattice system Λ and $A, B \subseteq \Lambda$ such that $A \cup B = \Lambda$.

PROBLEM

For a certain \mathcal{L}_Λ^* , can we prove $\alpha(\mathcal{L}_\Lambda^*) > 0$ using the result of quasi-factorization of the relative entropy?

THEOREM (C-Lucia-Pérez García, '18)

The **heat-bath dynamics**, with product fixed point, has a positive log-Sobolev constant.

Consider the local and global Lindbladians

$$\mathcal{L}_x^* := \mathbb{E}_x^* - \mathbb{1}_\Lambda, \quad \mathcal{L}_\Lambda^* = \sum_{x \in \Lambda} \mathcal{L}_x^*$$

Since

$$\mathbb{E}_x^*(\rho_\Lambda) = \sigma_\Lambda^{1/2} \sigma_{x^c}^{-1/2} \rho_{x^c} \sigma_{x^c}^{-1/2} \sigma_\Lambda^{1/2} = \sigma_x \otimes \rho_{x^c}$$

for every $\rho_\Lambda \in \mathcal{S}_\Lambda$, we have

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GENERAL QUASI-FACTORIZATION FOR σ A TENSOR PRODUCT

Let $\mathcal{H}_\Lambda = \bigotimes_{x \in \Lambda} \mathcal{H}_x$ and $\rho_\Lambda, \sigma_\Lambda \in \mathcal{S}_\Lambda$ such that $\sigma_\Lambda = \bigotimes_{x \in \Lambda} \sigma_x$. The following inequality holds:

$$D(\rho_\Lambda || \sigma_\Lambda) \leq \sum_{x \in \Lambda} D_x(\rho_\Lambda || \sigma_\Lambda).$$

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$$D(\rho_\Lambda || \sigma_\Lambda) \leq \sum_{x \in \Lambda} D_x(\rho_\Lambda || \sigma_\Lambda).$$

LOG-SOBOLEV CONSTANT

CONDITIONAL LOG-SOBOLEV CONSTANT

For $x \in \Lambda$, we define the **conditional log-Sobolev constant** of \mathcal{L}_Λ^* in x by

$$\alpha_\Lambda(\mathcal{L}_x^*) := \inf_{\rho_\Lambda \in \mathcal{S}_\Lambda} \frac{-\text{tr}[\mathcal{L}_x^*(\rho_\Lambda)(\log \rho_\Lambda - \log \sigma_\Lambda)]}{2D_x(\rho_\Lambda || \sigma_\Lambda)},$$

where σ_Λ is the fixed point of the evolution, and $D_x(\rho_\Lambda || \sigma_\Lambda)$ is the conditional relative entropy.

LEMMA

$$\alpha_\Lambda(\mathcal{L}_x^*) \geq \frac{1}{2}.$$

$$\begin{aligned}
 D(\rho_\Lambda || \sigma_\Lambda) &\leq \sum_{x \in \Lambda} D_x(\rho_\Lambda || \sigma_\Lambda) \\
 &\leq \sum_{x \in \Lambda} \frac{-\text{tr}[\mathcal{L}_x^*(\rho_\Lambda)(\log \rho_\Lambda - \log \sigma_\Lambda)]}{2\alpha_\Lambda(\mathcal{L}_x^*)} \\
 &\leq \frac{1}{2 \inf_{x \in \Lambda} \alpha_\Lambda(\mathcal{L}_x^*)} \sum_{x \in \Lambda} -\text{tr}[\mathcal{L}_x^*(\rho_\Lambda)(\log \rho_\Lambda - \log \sigma_\Lambda)] \\
 &= \frac{1}{2 \inf_{x \in \Lambda} \alpha_\Lambda(\mathcal{L}_x^*)} (-\text{tr}[\mathcal{L}_\Lambda^*(\rho_\Lambda)(\log \rho_\Lambda - \log \sigma_\Lambda)]) \\
 &\leq (-\text{tr}[\mathcal{L}_\Lambda^*(\rho_\Lambda)(\log \rho_\Lambda - \log \sigma_\Lambda)]) .
 \end{aligned}$$

POSITIVE LOG-SOBOLEV CONSTANT

$$\alpha(\mathcal{L}_\Lambda^*) \geq \frac{1}{2}.$$

EXAMPLES OF POSITIVE LOG-SOBOLEV CONSTANTS

ASSUMPTION 1

In a tripartite Hilbert space $\mathcal{H}_A \otimes \mathcal{H}_C \otimes \mathcal{H}_B$, A and B not connected, we have

$$\|h(\sigma_{AB})\|_\infty = \left\| \sigma_A^{-1/2} \otimes \sigma_B^{-1/2} \sigma_{AB} \sigma_A^{-1/2} \otimes \sigma_B^{-1/2} - \mathbf{1}_{AB} \right\|_\infty \leq K < \frac{1}{2}.$$

In particular, classical Gibbs states satisfy this.

ASSUMPTION 2

For any $B \subset \Lambda$, $B = B_1 \cup B_2$, it holds:

$$D_B(\rho_\Lambda || \sigma_\Lambda) \leq f(\sigma_{B\partial}) (D_{B_1}(\rho_\Lambda || \sigma_\Lambda) + D_{B_2}(\rho_\Lambda || \sigma_\Lambda)).$$

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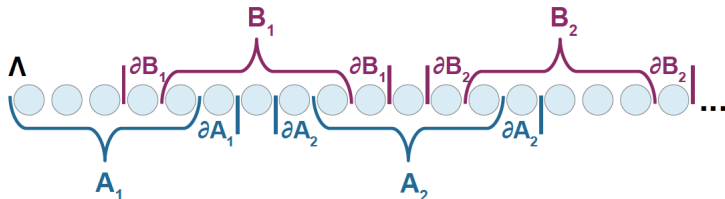
SKETCH OF THE PROOF

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STEP 1



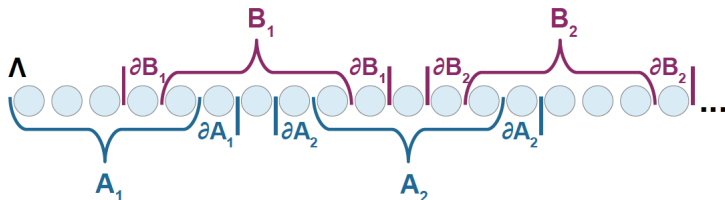
$$A = \bigcup_{i=1}^n A_i \text{ and } B = \bigcup_{j=1}^n B_j$$

$$D(\rho_\Lambda \| \sigma_\Lambda) \leq \frac{1}{1 - 2\|h(\sigma_{A^c B^c})\|_\infty} [D_A(\rho_\Lambda \| \sigma_\Lambda) + D_B(\rho_\Lambda \| \sigma_\Lambda)],$$

$$h(\sigma_{A^c B^c}) := \sigma_{A^c}^{-1/2} \otimes \sigma_{B^c}^{-1/2} \sigma_{A^c B^c} \sigma_{A^c}^{-1/2} \otimes \sigma_{B^c}^{-1/2} - \mathbb{1}_{A^c B^c}.$$

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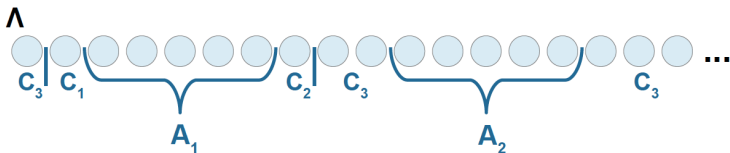


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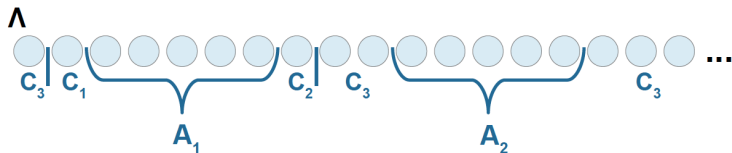


$$D_A(\rho_\Lambda || \sigma_\Lambda) \leq \sum_{i=1}^n D_{A_i}(\rho_\Lambda || \sigma_\Lambda)$$

σ_Λ is a QMC between $A_1 \leftrightarrow \partial A_1 \leftrightarrow \Lambda \setminus (A_1 \cup \partial A_1)$

$$\sigma_\Lambda = \bigoplus_{i \in I} \sigma_{A_1(\partial a_1)_i^L} \otimes \sigma_{(\partial a_1)_i^R \Lambda \setminus (A_1 \cup \partial A_1)}$$

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STEP 3

$$\text{Assumption 1} \Rightarrow \alpha(\mathcal{L}_\Lambda^*) \geq \tilde{K} \min_{i \in \{1, \dots, n\}} \{\alpha_\Lambda(\mathcal{L}_{A_i}^*), \alpha_\Lambda(\mathcal{L}_{B_i}^*)\}$$

STEP 4

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OPEN PROBLEMS

PROBLEM 1

Can we use any of the quasi-factorization results to prove log-Sobolev constants in a more general setting?

PROBLEM 2

Does this hold for greater dimension?

PROBLEM 3

Is there a better definition for conditional relative entropy?

